

**UNIVERSITY OF MARYLAND**  
**Department of Physics**  
**College Park, Maryland**

**PHYSICS Ph.D. QUALIFYING  
EXAMINATION PART II**

**September 3, 2020**

**11:00 a.m. – 3:00 p.m.**

**Do any four problems. Each problem is worth 25 points.  
Start each problem on a new sheet of paper (because different  
faculty members will be grading each problem in parallel).**

**Be sure to write your Qualifier ID (“control number”) at the top of  
each sheet — not your name! — and turn in solutions to four  
problems only. (If five solutions are turned in, we will only grade  
# 1 - # 4.)**

**At the end of the exam, when you are turning in your papers,  
please fill in a “no answer” placeholder form for the problem that  
you skipped, so that the grader for that problem will have  
something from every student.**

**You may keep this packet with the questions after the exam.**

### Problem II.1

Consider the problem of a heavy quark-antiquark bound state. Let the Hamiltonian be written, in the CM system, in a nonrelativistic form

$$H = \frac{\vec{p}^2}{2\mu} + V(\vec{r}), \quad (1)$$

where  $\vec{r} = \vec{r}_{quark} - \vec{r}_{antiquark}$  and  $\mu$  is the reduced mass of the quark and antiquark pair.

- (a) **[10 points]** Derive a relationship between the average kinetic energy and the average potential energy for the bound state. (Hint: Evaluate  $\frac{d}{dt}\langle\vec{r} \cdot \vec{p}\rangle$ ).
- (b) **[5 points]** The Feynman-Hellmann Theorem states that if the Hamiltonian  $H$  depends on a parameter  $\lambda$ ,  $H = H(\lambda)$ , and if  $|\psi(\lambda)\rangle$  is an eigenstate of  $H(\lambda)$ ,

$$H(\lambda)|\psi(\lambda)\rangle = E(\lambda)|\psi(\lambda)\rangle, \quad (2)$$

then

$$\frac{\partial}{\partial \lambda} E(\lambda) = \left\langle \psi(\lambda) \left| \left( \frac{\partial}{\partial \lambda} H(\lambda) \right) \right| \psi(\lambda) \right\rangle \quad (3)$$

Prove the theorem.

- (c) **[6 points]** Suppose

$$V(r) = V_0 \ln(r/r_0) \quad (4)$$

Deduce the dependence of the bound state energies  $E_n$  on the reduced mass  $\mu$ .

- (d) **[1 points]** From your results in (c) what, if anything, can you say about the dependence of  $E_n - E_m$  on the reduced mass  $\mu$ . (Here  $n$  and  $m$  refer generically to the quantum numbers of the bound states.)
- (e) **[3 points]** Estimate the radius of the ground state wavefunction, up to an unspecified factor of order unity.

## Problem II.2

Consider a quantum particle of mass  $m$  in a potential

$$V(x) = \frac{m\omega^2}{2}x^2 + cx^3. \quad (1)$$

We first consider the case with  $c = 0$ , and then add the  $x^3$  term using perturbation theory.

- (a) **[4 points]** What are the energy levels for the particle in the potential  $V(x)$  when  $c = 0$ ?
- (b) (i) **[4 points]** For what values of  $n$  and  $n'$  is the matrix element  $\langle n'|x^3|n\rangle$  between the unperturbed energy eigenstates nonzero? ( $n = 0$  labels the ground state.)  
(ii) **[4 points]** Write explicit expressions for the nonzero matrix elements  $\langle n'|x^3|n\rangle$ .
- (c) **[4 points]** What is the correction to the energy levels of Part (a) in first-order perturbation theory, where  $cx^3$  is the perturbation?
- (d) **[5 points]** Find the second-order correction to the ground state energy.
- (e) **[2 points]** For the potential  $V(x)$  of (1), the particle actually has no ground state. Explain why not, and describe the conditions under which it nevertheless makes physical sense to use perturbation theory for the ground state.
- (f) **[2 points]** What condition on  $c$  should be satisfied if second-order perturbation theory is to give a good approximation for the ground state?

You might find helpful the equation

$$\langle n'|x|n\rangle = \left(\frac{\hbar}{2m\omega}\right)^{1/2} \left(\sqrt{n}\delta_{n',n-1} + \sqrt{n+1}\delta_{n',n+1}\right)$$

where  $|n\rangle$  and  $|n'\rangle$  are orthogonal eigenstates of the unperturbed Hamiltonian.

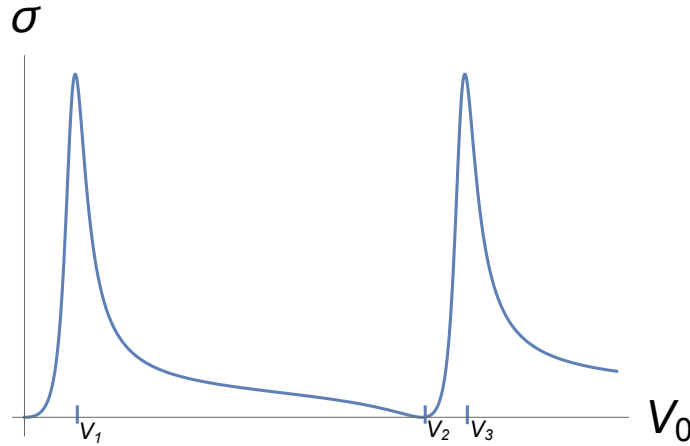
### Problem II.3

A particle of mass  $m$  and energy  $E = \hbar^2 k^2 / 2m$  scatters from a spherically symmetric potential,

$$V(r) = -V_0 \quad \text{for } r < R, \quad V(r) = 0 \quad \text{for } r > R. \quad (1)$$

In this problem we assume that the particle has very low energy  $E$ , such that  $kR \ll 1$ .

- (a) **[8 points]** Assuming that the strength  $V_0$  is not sufficient to support a bound state, find the cross section  $\sigma$  in the limit  $E \rightarrow 0$ , as a function of the remaining parameters,  $V_0$ ,  $R$ ,  $m$ , and  $\hbar$ .
- (b) **[4 points]** As the depth of the potential,  $V_0$ , is increased while keeping the energy  $E$  fixed, the total cross section behaves as in the Figure below.



Find the value of the total cross section at the maxima, where  $V_0 = V_1$  and  $V_0 = V_3$ , in terms of  $m$ ,  $E$ , and  $\hbar$ .

- (c) **[7 points]** Determine approximately the values  $V_1$  and  $V_3$  at which the first two maxima occur. *Hint:* Consider the values of  $V_1$  and  $V_3$  in the limit of zero energy.
- (d) **[6 points]** What is the approximate value of  $V_2$  at which the cross section is zero? The solution involves a transcendental equation. Sketch and label a graph showing the location of the relevant root, and give explicit upper and lower bounds for  $V_2$ .

Possibly useful formula connecting the amplitude  $f_l$  and the phase  $\delta_l$  in the scattering channel with the orbital angular momentum  $l$

$$f_l = \frac{e^{i\delta_l} \sin \delta_l}{k}.$$

### Problem II-4

Two particles interact via a spin-spin Hamiltonian  $A\mathbf{S}_1 \cdot \mathbf{S}_2$  where  $A$  is a positive constant and  $\mathbf{S}_{1,2}$  are the spin angular momenta of the two particles. Particle 1 has spin 1 and a magnetic moment of  $\mu_1 = -\frac{\mu_B}{\hbar}\mathbf{S}_1$ , whereas Particle 2 has spin  $\frac{1}{2}$  and magnetic moment zero.

- (a) **[6 points]** What are the energy levels of the system and the degree of degeneracy of the levels? Give a detailed derivation of the possible results.
- (b) **[8 points]** Write a basis of normalized energy eigenstates corresponding to the different energy levels in part (a) as linear combinations of products of single-particle spin states.
- (c) **[7 points]** If the system is placed in a magnetic field of strength  $B$  aligned with the z-axis, what then are the approximate energy eigenstates and eigenvalues if  $B \gg A\hbar^2/\mu_B$ ? Use the product of single particle spin states as before.
- (d) **[4 points]** Are any of the levels exactly linear in  $B$  for *all*  $B > 0$ ? If not, explain why not. If so, which levels are these and what are their energies as a function of  $B$ ?

*Possibly useful formula:*  $J_{\pm}|j, m\rangle = \sqrt{j(j+1) - m(m \pm 1)}|j, m \pm 1\rangle$ .

## Problem II.5

Consider a free-electron gas, with  $N$  electrons and dispersion relation  $\epsilon = (\hbar|\mathbf{k}|)^2/2m$  at temperature  $T=0$ , in  $d = 3$  or  $d = 2$  dimensions, contained in a volume  $V = L^3$  or area  $A = L^2$ , respectively, so a cube or a square.

- (a) **[6 points]** The number of single-particle energy eigenstates (counting all degeneracies) between  $\epsilon$  and  $\epsilon + d\epsilon$  is  $G(\epsilon) d\epsilon$ , where  $G(\epsilon)$  is known as the density of states. Show that  $G(\epsilon)$  satisfies

$$G(\epsilon) \propto \epsilon^\alpha, \quad (1)$$

finding the value of  $\alpha$  for  $d = 3$  and for  $d = 2$ . Assume throughout that  $L$  is sufficiently large that finite-size effects can be ignored. Ignore numerical factors and dimensionful constants, since you seek only the  $\epsilon$  dependence of  $G(\epsilon)$ . (*Hint:* With periodic boundary conditions each single-particle state can be taken to have a definite wave vector.)

- (b) **[8 points]** i) (1) What are the units of  $G(\epsilon)$ ?  
(2) Explain why the proportionality relation (1) can be written more specifically as the following equation:

$$G(\epsilon) = \frac{BN\epsilon^\alpha}{\epsilon_F^{\alpha'}}, \quad (2)$$

where  $\epsilon_F$  is the Fermi energy and  $B$  is a numerical constant. Specify how the numerical value of  $\alpha'$  is related to that of  $\alpha$ . (*Note:* To answer this question one does not need to determine  $\epsilon_F$  in terms of the parameters of the problem.)

ii) Find the value of  $B$  for  $d = 3$  and for  $d = 2$ .

iii) For  $N$  electrons, how does  $\epsilon_F$  depend on  $N$  and  $V$ , and on  $N$  and  $A$  for  $d = 2$ ? [Again, prefactors are not needed, just the proportionality with the correct exponents.]

For graphene (a single planar sheet of graphite) near the Dirac points, the electronic dispersion relation can be written  $\epsilon = \hbar v_s |\mathbf{k}|$  (where  $\mathbf{k}$  is measured from a Dirac point, a fact you can ignore here).

- (c) **[6 points]** i) Show that the density of states  $G(\epsilon)$  is still proportional to  $\epsilon^\alpha/\epsilon_F^{\alpha'}$ , and find the new value of  $\alpha$  for  $d = 2$ .  
ii) Does the relationship of  $\alpha'$  to  $\alpha$  change from part 2? If yes, how? If not, why not?
- (d) **[5 points]** The total low-temperature heat capacity of a metal with fixed volume  $V$  [or area  $A$ ] and fixed  $N$  is known to behave as a power-law of  $T$ . Give a quick argument to show what this power is. It may be helpful to sketch the change in the Fermi-Dirac distribution when  $T$  increases slightly from 0. (Use of the Sommerfeld expansion is not intended!)